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Course 565/665 Lecturer prof. Cavagnero
 Day 4.2.04 Date 9:55am
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predict μ_A, μ_B from microscopic model.

for Ideal Gas. $q' = \sum_{j=0}^{\infty} e^{-\beta E_j} = e^{-\epsilon_0/kT} + e^{-\epsilon_1/kT} + \dots$

define $q = q' \cdot e^{\epsilon_0/kT} = 1 + e^{-(\epsilon_1 - \epsilon_0)/kT} + e^{-(\epsilon_2 - \epsilon_0)/kT} + \dots$

for IG. can show: $\mu = -kT \ln\left(\frac{q}{N}\right)$

$$\left\{ \begin{array}{l} \mu_A = -kT \ln \frac{q'_A}{N_A} \\ \mu_B = -kT \ln \frac{q'_B}{N_B} \end{array} \right. \xrightarrow{\text{at eq.}} \frac{q'_B}{N_B} = \frac{q'_A}{N_A}$$

$$K_{AB} = \frac{N_B}{N_A} = \frac{q'_B}{q'_A} = \frac{q_B}{q_A} \cdot e^{-(\epsilon_{0B} - \epsilon_{0A})/kT}$$

In general: $aA + bB \rightleftharpoons cC$

same approach leads to: $K = \frac{q'_C}{q'_A{}^a q'_B{}^b} e^{-(c\epsilon_{0C} - a\epsilon_{0A} - b\epsilon_{0B})/kT}$

show $\mu_A = -kT \ln \frac{q}{N}$ for IG.

$$Q = \frac{q^N}{N!} \quad (\text{independent, indistinguishable})$$

$$\begin{aligned} G &= -kT \ln Q = -kT \ln \frac{q^N}{N!} = -kT [\ln q^N - \ln(N!)] \\ &= -kT \ln \left(\frac{qe}{N} \right)^N \end{aligned}$$

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$$\mu = \frac{\partial}{\partial N} \left[kT N \ln \left(\frac{N}{q_e} \right) \right] = kT \ln \left(\frac{N}{q_e} \right) + kT N \frac{1}{N} \cdot \frac{1}{q_e}$$
$$= -kT \ln \frac{q}{N}$$

see. $K_{AB} = \frac{q_B}{q_A} \cdot e^{-(\epsilon_{0B} - \epsilon_{0A})/kT}$

if $q_A \sim q_B$, if $\epsilon_{0A} < \epsilon_{0B} \Rightarrow K_{AB}$ is small.

if $\epsilon_{0A} < \epsilon_{0B}$, but $q_A < q_B$, could have $K_{AB} > 1$.

K_{AB} : the balance between ratio of partition fn. and difference of ground state energy.

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