Relaxation and the Nuclear Overhauser Effect

- Relaxation rates depends on matching transition frequencies to timevarying magnetic fields.
- A proton affects surrounding protons via dipole-dipole interactions. The dipole field can be visualized as a small bar magnet placed at the proton nucleus.

$$\Delta E_{dd} \propto \gamma_i^2 \gamma_j^2 \left\langle \sum_j \frac{\tau_c}{r_{ij}^6} \right\rangle_t$$



• Note the distance and correlation time dependencies.

Molecular Motions and Relaxation

 Relaxation rates depend on matching transition frequencies to timevarying magnetic fields: T₁ relaxation is most efficient when protons experience molecular motions at the Larmor frequency.



The Nuclear Overhauser Effect (2-spin system)

In NMR: excess population \propto 1 - exp(- Δ E/RT) ~ Δ E

(see Claridge or Sanders&Hunter)



The Nuclear Overhauser Effect

• Through-space dipole-dipole interactions!

[spins can be J-coupled, but nOe does not arise from J-coupling!]



The Nuclear Overhauser Effect

• Start with selective inversion:



NOEs in the Rotating Frame: NOEs

• Selective inversion can be pictured in a simplified manner as:



Molecular Motions and Relaxation

 Relaxation rates depend on matching transition frequencies to timevarying magnetic fields: T₁ relaxation is most efficient when protons experience molecular motions at the Larmor frequency.



Negative NOEs: Large MW

• Slow motions are effective for small frequencies/energy differences:

 $W_0(ZQ) \Rightarrow$ negative NOE dominates for large MW



Molecular Motions and Relaxation

 Relaxation rates depend on matching transition frequencies to timevarying magnetic fields: T₁ relaxation is most efficient when protons experience molecular motions at the Larmor frequency.



Positive NOEs: Small MW

Fast motions are effective for large frequencies/energy differences:
W₂ (DQ) ⇒ positive NOE dominates for small MW



Zero NOEs: The Crossover Region

 Thus, NOEs go through a *crossover region* at intermediate MW, in the range 1000-5000 Da.



The ROESY Spin-Lock: Reducing the Effective Field

 Problems with the crossover region can be avoided by spin-locking the magnetization.

During a spin-lock, the *effective* magnetic field (on-resonance) is B_1 .



Claridge

Figure 5.62 A schematic illustration of events during spin-lock mixing. All chemical shift differences between spins are eliminated yet all spin–spin couplings between them remain. This forces the *strong-coupling* condition on all spins (see text).



"All" molecules in

solution tumble fast

compared to 6 kHz!

ROEs are always positive.

Review - ROESY

- Theoretical steady-state ROESY enhancements are shown below.
 - ⇒ In practice, the spin-lock causes many problems, the worst being the possibility of TOCSY (J-coupling) transfers in the spectrum.
 - × Avoid having coupled multiplets centered in the spectrum.
 - ⇒ A number of variations of ROESY exist, with differing attributes.
 - ⇒ Attempt NOESY first, and use ROESY only if required.



Heteronuclear NOEs: Positive γ Nuclei



Heteronuclear NOEs: Positive γ Nuclei







Homonuclear NOEs: Transient Experiments

Transient NOE experiments (e.g., NOESY1D, NOESY2D) will impose a limitation on the mix time of approx T_1 . At mix ~ T_1 and longer, loss of magnetization via spin-lattice relaxation will dominate over the NOE.

Larger %NOE is an indication of smaller r, *but definitely not always*: spin diffusion causes problems for large MW, and certain geometries are problematic for all MW.

η



Homonuclear NOEs: Transient Experiments

Transient NOE experiments (e.g., NOESY1D, NOESY2D) will impose a limitation on the mix time of approx T_1 . At mix ~ T_1 and longer, loss of magnetization via spin-lattice relaxation will dominate over the NOE.

Larger %NOE is an indication of smaller r, <u>but definitely not</u> <u>always</u>: spin diffusion causes problems for large MW, and certain geometries are problematic for all MW. And T_1 !

η



NOE - Qualitative and Quant Cautions

- Geometry can be very important to NOE interpretations.
 - \Rightarrow E.g., three protons distributed in a near-equilateral triangle can produce zero NOE, independent of r_{IS} .
 - ⇒ It is important to be aware that *not* observing an NOE is weak evidence. *Measure NOEs in all directions.*



• The enhancement maximizes at $1 + \gamma_H/2\gamma_X$ (fast limit) or 0 (slow limit).

Note lack of r dependence!!



Claridge Figure 8.12

• The enhancement maximizes at $1 + \gamma_H/2\gamma_X$ (fast limit) or 0 (slow limit).

Note lack of r dependence!!



• The build-up rate with mix time is r dependent:

 $R_v = K' \gamma_I^2 \gamma_S^2 \tau_c r_{IS}^{-6}$

Plot NOE versus mix, and compare to known pair:



• The build-up rate with mix time is r dependent:

 $R_v = K' \gamma_I^2 \gamma_S^2 \tau_c r_{IS}^{-6}$

Plot NOE versus mix, and compare to known pair:



Summary: NOE

- NOE's occur via population transfers, and are slow to occur (taking times approaching ~T₁).
 - DQ relaxation leads to positive NOEs, but requires high-frequency modulations (low MW).
 - ZQ relaxation leads to negative NOEs, occurs for all molecules (0 frequency): dominates at high MW.
- NOESY-1D and NOESY-2D are *transient* experiments. These experiments have utility for mix ≤ T₁.
- The spin-lock of a ROESY experiment reduces the effective magnetic field to 2-6 kHz. Thus, ROEs are *always* positive (i.e., motions in liquids are *always* fast compared to these kHz frequencies).

Review - TOCSY and Scalar Coupling

• **Spin-locking** the magnetization scales the chemical shift to nearzero (in Hz), producing strong coupling.



Review - TOCSY and Scalar Coupling

• **Spin-locking** the magnetization scales the chemical shift to nearzero (in Hz), producing strong coupling.



Review - TOCSY and Scalar Coupling

• **Spin-locking** the magnetization scales the chemical shift to nearzero (in Hz), producing strong coupling.

Strongly coupled protons exchange magnetization once every ~1/2J.



Claridge

Figure 5.62 A schematic illustration of events during spin-lock mixing. All chemical shift differences between spins are eliminated yet all spin–spin couplings between them remain. This forces the *strong-coupling* condition on all spins (see text).

